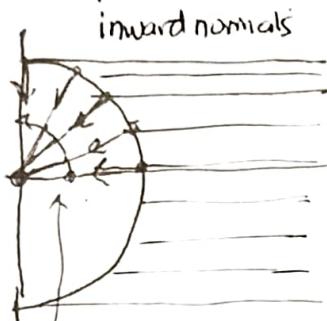
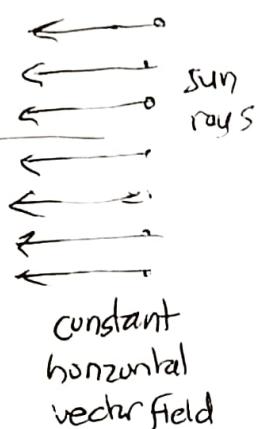


Flux in 2-d

consider the half of the Earth facing the sun — solid of revolution about symmetry axis connecting the centers of the Earth and sun
consider a plane cross-section through that symmetry axis

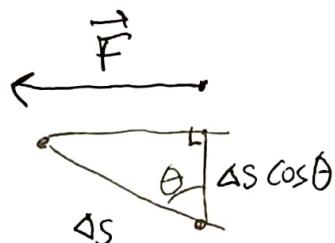
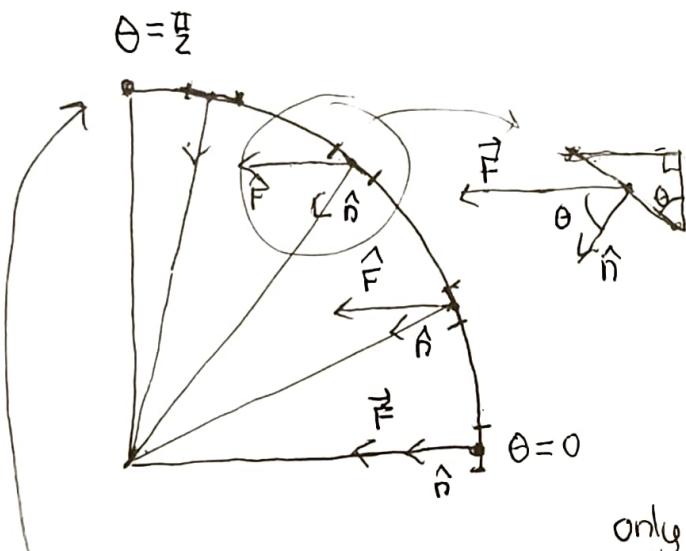


θ is angle with horizontal radius R



\vec{F} is flux of energy hitting a line perpendicular to the rays per unit length

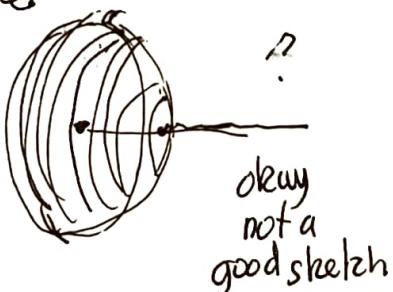
as we approach $\theta = \pi/2$,
equally spaced increments in θ
(hence of $\Delta S = R\Delta\theta$)
receive less and less energy
because of the "surface" tilting
with respect to the rays



$$|\vec{F}|(\Delta S \cos\theta) = \vec{F} \cdot \hat{n} \Delta S = F_{\perp} \Delta S$$

only the normal component of the impinging light rays is effective in delivering energy to the surface — at the tips the rays are parallel & so deliver no energy: $\cos\theta \rightarrow 0$ there

Of course this reduces around the axis to a 3-d problem of annular rings on the Earth's surface.



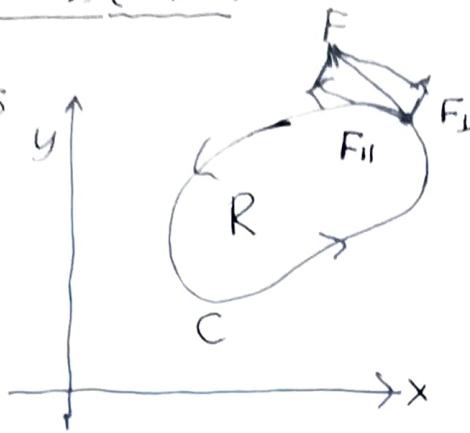
energy per unit area perpendicular to sun rays

$|\vec{F}|(\Delta A)_{\perp}$ is amount of energy delivered to any amount of surface area

integrating up the normal component (2-d or 3-d) against the arclength/area of the curve/surface measures the "flux" of that vector field through that curve/surface.

Green-Stokes-Gauss (8 flux)

2-d Greens



$$F_{\perp} = \hat{n} \cdot \vec{F}$$

$$F_{\parallel} = \hat{T} \cdot \vec{F}$$

\hat{n} = outward normal

\hat{T} = counterclockwise

unit normal

$$\hat{n} \times \hat{T} = \hat{k}$$

righthand rule!

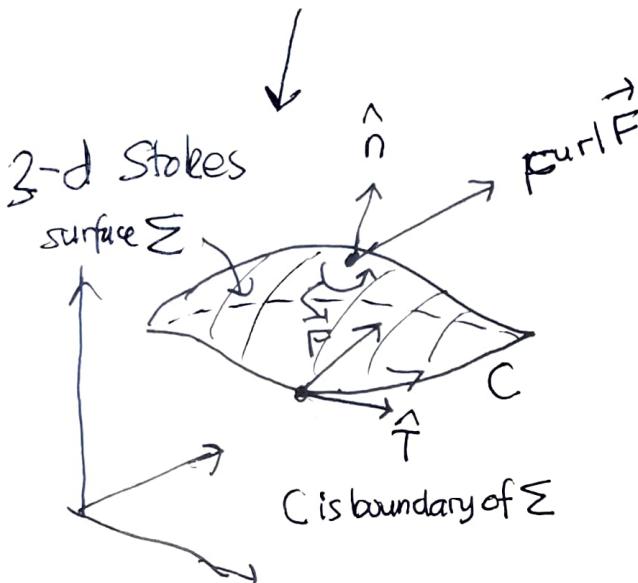
Stokes version:

$$\oint_C \vec{F} \cdot d\vec{r} = \oint_C F_{\parallel} ds$$

circulation of \vec{F}
around C

$$= \iint_R (\operatorname{curl} \vec{F}) \cdot \hat{k} dA$$

$\underbrace{R}_{\text{"circulation density"}}$



$$\oint_C \vec{F} \cdot d\vec{r} = \oint_C F_{\parallel} ds$$

$$= \iint_{\Sigma} (\operatorname{curl} \vec{F}) \cdot \hat{n} dS$$

(flux of $\operatorname{curl} \vec{F}$
thru Σ in
direction \hat{n})

dS vector differential
of surface area

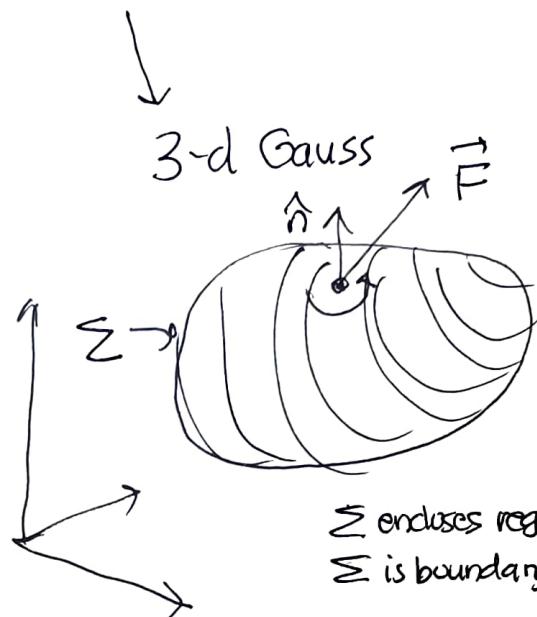
Gauss version:

$$\oint_C F_{\perp} ds$$

"flux" of \vec{F}
out of R
thru C

$$= \iint_R (\operatorname{div} \vec{F}) dA$$

$\underbrace{R}_{\text{"flux density"}}$



$$\iint_{\Sigma} \vec{F} \cdot d\vec{S} = \iint_{\Sigma} F_{\perp} dS$$

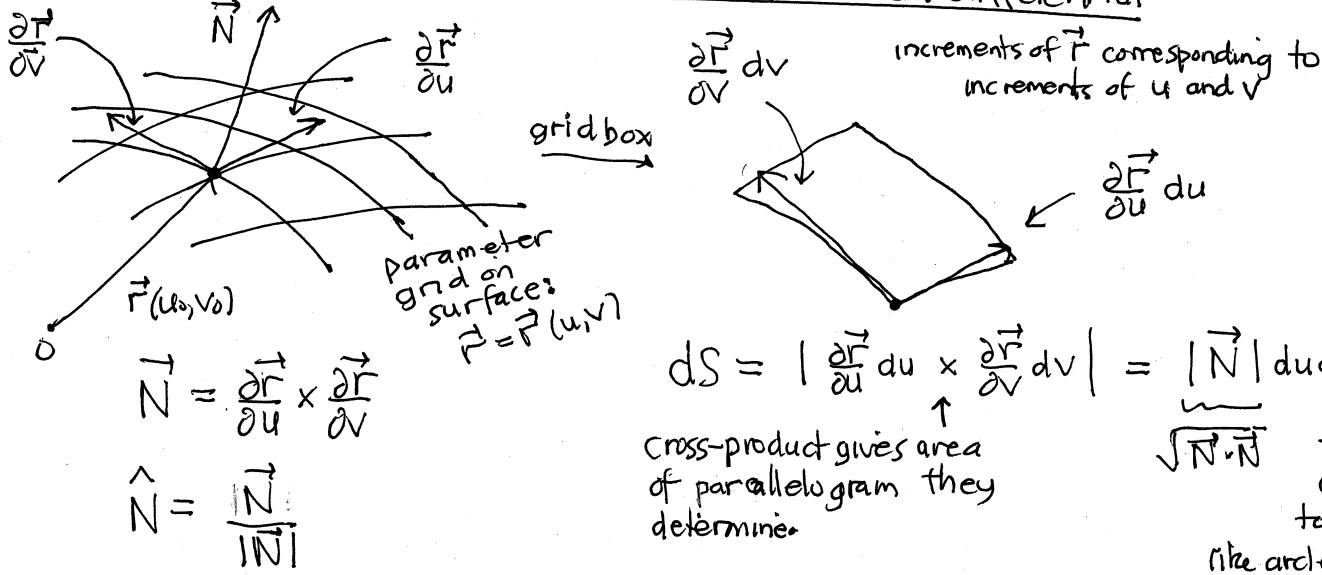
(flux of
 \vec{F} thru
 Σ in
direction \hat{n})

$$= \iiint_R (\operatorname{div} \vec{F}) dV$$

in both cases:

integral of vector field on boundary = integral of its appropriate derivative
on the interior region inside the boundary

surface integral: normal vector and surface area differential



surface integrals integrate the normal component of a vector field against the differential of surface area.

$$\iint_{\text{surface}} \vec{F}(\vec{r}(u, v)) \cdot \hat{N}(\vec{r}(u, v)) \, dS$$

$\underbrace{|\vec{N}(\vec{r}(u, v))|}_{\vec{N}(\vec{r}(u, v))} \, dudv$

$$= \iint_{U_1 U_2} \vec{F}(\vec{r}(u, v)) \cdot \vec{N}(\vec{r}(u, v)) \, dv \, du$$

square root goes away
for vector integrals

easier than integrating scalars

$$\left[\iint_C \vec{F}(\vec{r}(t)) \cdot \hat{T}(\vec{r}(t)) \, dS \right]$$

$\underbrace{|\vec{F}(\vec{r}(t))|}_{\vec{F}'(t)} \, dt$

$$= \int_{t_1}^{t_2} \vec{F}(\vec{r}(t)) \cdot \vec{F}'(t) \, dt$$

just like for the
line integral case

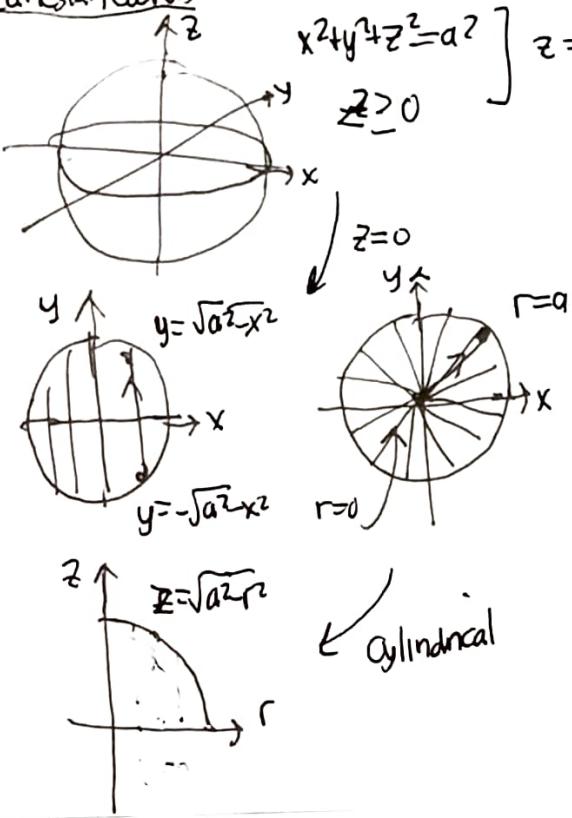
Scalar surface integrals

Like scalar line integrals, sart factor makes them more difficult to evaluate

$$\iint_S f dS = \int_{U_1}^{U_2} \int_{V_1}^{V_2} f(\vec{r}(u, v)) \underbrace{\|\vec{N}(\vec{r}(u, v))\| du dv}_{\left| \frac{\partial \vec{r}}{\partial u} \times \frac{\partial \vec{r}}{\partial v} \right|} = \sqrt{\dots} \text{ sart factor}$$

$$\text{surface area of hemisphere} = 2\pi a^2 = \frac{1}{2}(4\pi a^2) = S$$

cartesian coords



cylindrical

spherical coords

$$\vec{r} = a \langle \sin\phi \cos\theta, \sin\phi \sin\theta, \cos\phi \rangle = a \hat{r}$$

$$\frac{\partial \vec{r}}{\partial \phi} = a \langle \cos\phi \cos\theta, \cos\phi \sin\theta, -\sin\phi \rangle$$

$$\frac{\partial \vec{r}}{\partial \theta} = a \langle -\sin\phi \sin\theta, \sin\phi \cos\theta, 0 \rangle$$

$$\frac{\partial \vec{r}}{\partial \phi} \times \frac{\partial \vec{r}}{\partial \theta} = \dots = a^2 \sin\phi \langle \sin\phi \cos\theta, \sin\phi \sin\theta, \cos\phi \rangle = a^2 \sin\phi \hat{r} \quad (\text{radial outward})$$

$$\left| \frac{\partial \vec{r}}{\partial \phi} \times \frac{\partial \vec{r}}{\partial \theta} \right| = a^2 \sin\phi$$

$$S = \int_0^{2\pi} \int_0^{\pi/2} a^2 \sin\phi \, d\phi \, d\theta = 2\pi a^2 \int_0^{\pi/2} \sin\phi \, d\phi = 2\pi a^2 \left[-\cos\phi \right]_0^{\pi/2} = 2\pi a^2$$

$$\frac{\partial \vec{r}}{\partial x} = \langle 1, 0, \frac{1}{z} \sqrt{a^2 - x^2 - y^2} \rangle = \langle 1, 0, -\frac{x}{z} \rangle$$

$$\frac{\partial \vec{r}}{\partial y} = \langle 0, 1, -\frac{y}{z} \rangle$$

$$\frac{\partial \vec{r}}{\partial x} \times \frac{\partial \vec{r}}{\partial y} = \dots = \langle \frac{x}{z}, \frac{y}{z}, 1 \rangle = \frac{\langle x, y, z \rangle}{z} = \frac{\vec{r}}{z}$$

radially outward!

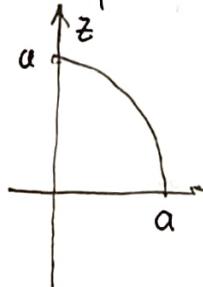
$$\left| \frac{\partial \vec{r}}{\partial x} \times \frac{\partial \vec{r}}{\partial y} \right| = \frac{\sqrt{x^2 + y^2 + z^2}}{z} = \frac{a}{z} = \frac{a}{\sqrt{a^2 - x^2 - y^2}}$$

$$S = \underbrace{\int_{-a}^a \int_{-\sqrt{a^2 - x^2}}^{\sqrt{a^2 - x^2}}}_{\int_0^{2\pi} \int_0^a} dy \, dx \quad \text{hard}$$

$$\underbrace{\frac{a}{\sqrt{a^2 - x^2}}}_{\text{u-sub!}} \underbrace{r dr d\theta}_{\text{easier}} \quad = 2\pi a^2 \checkmark$$

Scalar surface integrals (2)

nemisphere: cylindrical coords



$$r^2 + z^2 = a^2$$

$$z = \sqrt{a^2 - r^2}$$

$$\vec{r} = \langle r\cos\theta, r\sin\theta, \sqrt{a^2 - r^2} \rangle, \quad r=0..a, \quad \theta=0..2\pi$$

$$\frac{\partial \vec{r}}{\partial r} = \langle \cos\theta, \sin\theta, \frac{(-2r)}{2\sqrt{a^2 - r^2}} \rangle = \langle \cos\theta, \sin\theta, -r/z \rangle$$

$$\frac{\partial \vec{r}}{\partial \theta} = \langle -r\sin\theta, r\cos\theta, 0 \rangle \equiv \langle -rs, rc, 0 \rangle$$

"radial": $\frac{\partial \vec{r}}{\partial r} \times \frac{\partial \vec{r}}{\partial \theta} = \dots = r \left\langle \frac{r}{z} c, \frac{r}{z} s, 1 \right\rangle = \frac{r}{z} \underbrace{\langle r\cos\theta, r\sin\theta, z \rangle}_{\vec{r}}$

$$\left| \frac{\partial \vec{r}}{\partial r} \times \frac{\partial \vec{r}}{\partial \theta} \right| = r \sqrt{\frac{r^2}{a^2 - r^2} (\underbrace{c^2 + s^2}_1) + 1} = r \sqrt{\frac{r^2 + (a^2 - r^2)}{a^2 - r^2}} = \underbrace{\left(\frac{a}{\sqrt{a^2 - r^2}} \right)}_{\text{comes from } (\vec{N})} r$$

$$dS = \underbrace{\left(\frac{a}{\sqrt{a^2 - r^2}} \right)}_{r} dr d\theta$$

comes from (\vec{N}) , automatically

↓ same as cartesian integral re-expressed in cylindrical coords.