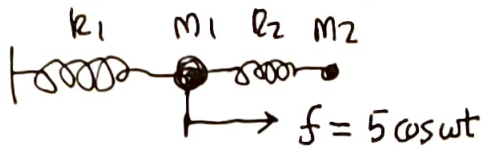


EPC4 7.5.14 dynamic damping

①



$k_1 = 50, k_2 = 10, m_1 = 1, \omega = 10$

$$\begin{bmatrix} x_1 \\ x_2 \end{bmatrix}'' = \underbrace{\begin{bmatrix} -(k_1+k_2)/m_1 & k_2/m_1 \\ k_2/m_2 & -k_2/m_2 \end{bmatrix}}_A \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} + \underbrace{\begin{bmatrix} 5 \cos wt \\ 0 \end{bmatrix}}_F$$

$$A = \begin{bmatrix} -60 & 10 \\ 10/m_2 & -10/m_2 \end{bmatrix}$$

Goal: Adjust m_2 to make the response displacement for x_1 zero. Use the method of undetermined coefficients for x_1, x_2 to evaluate the particular soln, which is the response vector.

Trial soln: $x = \langle x_1, x_2 \rangle = \langle a \cos wt, b \cos wt \rangle$

$$x'' = -\omega^2 x = Ax + F$$

$$\underbrace{(A + \omega^2 I)}_{} x = F \rightarrow x = \underbrace{-(A + \omega^2 I)^{-1}}_{\text{insert}} F$$

$\det(A - (-\omega^2)I) \neq 0$
 \neq eigenvalue of A

coefficients of $\cos wt$:

Explicitly:

$$x_1'' + 60x_1 - 10x_2 = 5 \cos 10t \rightarrow -100a + 60a - 10b = 5$$

$$x_2'' - \frac{10}{m_2}x_1 + \frac{10}{m_2}x_2 = 0 \rightarrow -100b - \frac{10}{m_2}a + \frac{10}{m_2}b = 0$$

$$\begin{bmatrix} -40 & -10 \\ \frac{10}{m_2} & \frac{10}{m_2} - 100 \end{bmatrix} \begin{bmatrix} a \\ b \end{bmatrix} = \begin{bmatrix} 5 \\ 0 \end{bmatrix} \rightarrow \begin{bmatrix} a \\ b \end{bmatrix} = \frac{\begin{bmatrix} \frac{10}{m_2} - 100 & 10 \\ \frac{10}{m_2} & -40 \end{bmatrix} \begin{bmatrix} 5 \\ 0 \end{bmatrix}}{500(8 - 1/m_2)} = \begin{bmatrix} \frac{50(\frac{1}{m_2} - 10)}{500(8 - 1/m_2)} \\ \frac{50/m_2}{500(8 - 1/m_2)} \end{bmatrix}$$

$\det: 40(100 - \frac{10}{m_2}) - \frac{100}{m_2} = 500(8 - \frac{1}{m_2})$

if $m_2 = \frac{1}{8}$, $\det(A + \omega^2 I) = 0$ and $\lambda = -100$ is an eigenvalue corresponding to eigenfrequency $\omega = 10$.

In this case the trial solution must be multiplied by t because the trial function is a homogeneous soln!

$$= \begin{bmatrix} \frac{1}{m_2} - 10 \\ \frac{1}{m_2} \end{bmatrix} \frac{10(8 - 1/m_2)}{10(8 - 1/m_2)} \quad \begin{matrix} a=0; m_2 = 1/10 \\ \rightarrow \\ b = \frac{100}{10(8-10)} = -\frac{1}{2} \end{matrix}$$

$$x = \langle 0, -\frac{1}{2} \cos 10t \rangle$$

↑ ↑
 mass 1 is fixed, mass 2 oscillates in opposite direction to the force on mass 1 to compensate for the force on it from the second spring.

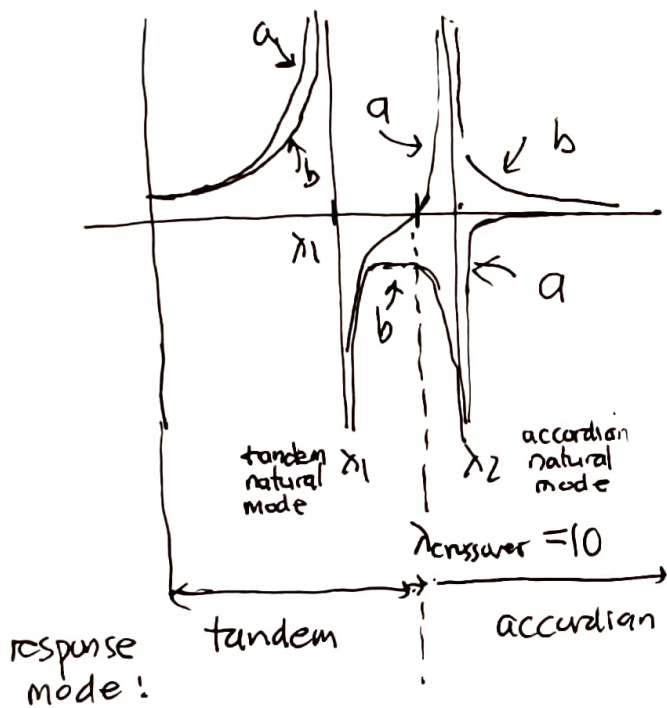
EPC4 7.6.14 dynamic damping (2)

For $m = 1/8$:

$$A = \begin{bmatrix} -60 & 10 \\ 100 & -100 \end{bmatrix} \quad B \approx \begin{bmatrix} 0.57 & -0.17 \\ 1 & 1 \end{bmatrix}$$

$$\lambda \approx \begin{matrix} 6.53 & 10.84 \\ \omega_1 & \omega_2 \\ \text{tandem} & \text{accordian} \\ \text{mode} & \text{mode} \end{matrix}$$

If we study the response coefficients $[a(\omega), b(\omega)]$ for general frequency we find



This is typical for all of these 2 mass 2 spring systems.

The interval $0 \leq \omega \leq \omega_{\text{crossover}}$ containing the tandem natural mode frequency has a tandem mode response (same sign for $a(\omega), b(\omega)$)

for $\omega > \omega_{\text{crossover}}$ containing the accordian mode frequency, the response mode is a accordian mode (opposite signs for $a(\omega), b(\omega)$)

The first mass is stationary at the crossover frequency, where the second spring contracts or stretches to balance the force applied directly to the first mass.

$$\begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = \begin{bmatrix} a(\omega) \cos \omega t \\ b(\omega) \cos \omega t \end{bmatrix}$$

$$f_1 = 5 \cos \omega t$$

↑
just a convenient factor
actual value unimportant.

It is the ratio of the applied force and the response amplitudes which matter.

For fixed M_1, k_1 one can always adjust the smaller mass-spring add on to zero out the oscillations in the first mass.

EPC4 7.6.14 dynamic damping (more!)

(3)

$$A = \begin{bmatrix} \frac{-(k_1+k_2)}{m_1} & \frac{k_2}{m_1} \\ \frac{k_2}{m_2} & -\frac{k_2}{m_2} \end{bmatrix} = \begin{bmatrix} -(\omega_1^2 + \gamma\omega_2^2) & \omega_2^2 \gamma \\ \omega_2^2 & -\omega_2^2 \end{bmatrix}$$

$$\omega_1^2 \equiv \frac{k_1}{m_1}$$

$$\omega_2^2 \equiv \frac{k_2}{m_2}$$

separate natural frequencies

$$\gamma = \frac{m_2}{m_1} \text{ mass ratio}$$

$$x = \begin{bmatrix} a \\ b \end{bmatrix} \cos \omega t, \quad x'' = -\omega^2 x$$

$$\downarrow$$

$$x'' - Ax = \begin{bmatrix} \omega_1^2 + \gamma\omega_2^2 - \omega^2 & -\gamma\omega_2^2 \\ -\omega_2^2 & \omega_2^2 - \omega^2 \end{bmatrix} \begin{bmatrix} a \\ b \end{bmatrix} \cos \omega t = \begin{bmatrix} f_1 \cos \omega t \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} \omega_1^2 + \gamma\omega_2^2 - \omega^2 & \gamma\omega_2^2 \\ -\omega_2^2 & \omega_2^2 - \omega^2 \end{bmatrix} \begin{bmatrix} a \\ b \end{bmatrix} = \begin{bmatrix} f_1 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} a \\ b \end{bmatrix} = \frac{1}{(\omega_1^2 + \gamma\omega_2^2 - \omega^2)(\omega_2^2 - \omega^2) + \gamma\omega_2^4} \begin{bmatrix} \omega_2^2 - \omega^2 & -\gamma\omega_2^2 \\ \omega_2^2 & \omega_1^2 + \gamma\omega_2^2 - \omega^2 \end{bmatrix} \begin{bmatrix} f_1 \\ 0 \end{bmatrix} = \frac{f_1}{\mathcal{D}} \begin{bmatrix} \omega_2^2 - \omega^2 \\ \omega_2^2 \end{bmatrix}$$

$$\mathcal{D} \rightarrow a=0 \text{ means } \omega^2 = \omega_2^2, \quad \boxed{\omega = \omega_2}$$

The second mass natural frequency is used to drive the first mass or more precisely

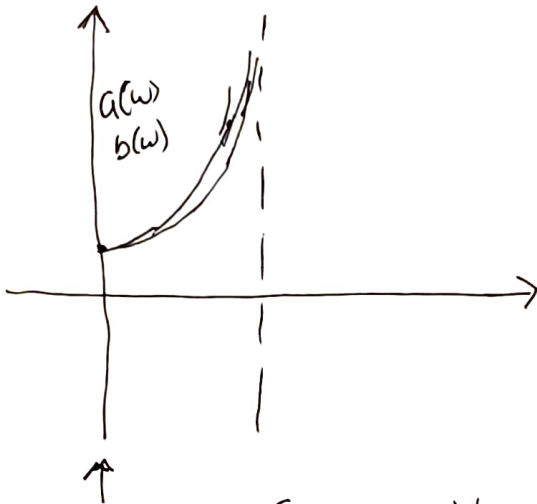
The second mass was chosen to make the second mass spring natural frequency equal the driving frequency, which in turn is slightly less than the natural frequency of the combined system (accadian mode) where resonance occurs (when some damping is present).

Pushing this concrete problem with specific numbers to understand what lies behind it is the magic of mathematics.

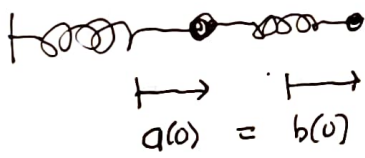
But we can actually understand some obvious features of the frequency response functions $a(\omega), b(\omega)$.

EPC4 7.6.14 dynamic damping

4



$\omega=0$ zero frequency limit
 $\cos(0) = 1$: constant force
on first mass to right
both masses just move to
new equilibrium equal distances
to the right



since there is no third spring
to compress the second spring

As ω increases from zero
to very small values, the system
simply slowly moves to new
equilibrium values moment by
moment so remains in phase
with the slowly changing
driving force.